

How to distinguish two
mass-degenerate 125 GeV
Higgs boson(s) from one?

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Outline

- Can we have more than one (or two) resonances at 125 GeV (degenerate Higgs)? Yes, it is possible.
- How to distinguish the degenerate Higgs scenario from the conventional single Higgs scenario?
 - precise signal fit cannot offer an answer
 - to look for the BSM effects: self-interactions or pair productions at LHC/ILC, the searches for CP-odd scalars.

The 2HDM

$$\mathcal{L} = \sum_{i=1,2} |D\Phi_i|^2 - V(\Phi_1, \Phi_2)$$

$$V(\Phi_1, \Phi_2) = m_{11}^2 |\Phi_1|^2 + m_{22}^2 |\Phi_2|^2 - (m_{12}^2 \Phi_1^\dagger \Phi_2 + H.c.) + \frac{1}{2} \lambda_1 |\Phi_1|^4 + \frac{1}{2} \lambda_2 |\Phi_2|^4 \\ + \lambda_3 |\Phi_1|^2 |\Phi_2|^2 + \lambda_4 |\Phi_1^\dagger \Phi_2|^2 + \frac{1}{2} [\lambda_5 (\Phi_1^\dagger \Phi_2) (\Phi_1^\dagger \Phi_2) + H.c.]$$

with m_{12}^2 and λ_5 being complex.

$$\Phi_1 = \begin{pmatrix} -s_\beta H^+ \\ \frac{1}{\sqrt{2}} (v_1 + H_1^0 - i s_\beta A) \end{pmatrix}, \quad \Phi_2 = \begin{pmatrix} c_\beta H^+ \\ \frac{1}{\sqrt{2}} (v_2 e^{i\xi} + H_2^0 + i c_\beta A) \end{pmatrix}$$

Neutral Higgs masses

- Masses and mixings of the CP-even Higgs boson:

$$\mathcal{M}_0^2 = \begin{pmatrix} m_{12}^2 t_\beta + \lambda_1 v_1^2 & -m_{12}^2 + \lambda_{345} v_1 v_2 \\ -m_{12}^2 + \lambda_{345} v_1 v_2 & m_{12}^2 / t_\beta + \lambda_2 v_2^2 \end{pmatrix},$$
$$\begin{pmatrix} H \\ h \end{pmatrix} = \begin{pmatrix} c_\alpha & s_\alpha \\ -s_\alpha & c_\alpha \end{pmatrix} \cdot \begin{pmatrix} h_1 \\ h_2 \end{pmatrix}$$

Neutral Higgs couplings

$$\mathcal{L} \supset \sum_{h_i=h,H} \left[-\frac{m_f}{v} \xi_i^f \bar{f} f + a_i \left(2 \frac{m_W^2}{v} W_\mu^+ W^{-\mu} + \frac{m_Z^2}{v} Z_\mu Z^\mu \right) \right] h_i - \frac{m_f}{v} \xi_A^f \bar{f} i \gamma_5 f A,$$

$$\text{Type-I: } \xi_h^f = s_{\beta-\alpha} + \frac{c_{\beta-\alpha}}{t_\beta}, \quad \xi_H^f = c_{\beta-\alpha} - \frac{s_{\beta-\alpha}}{t_\beta}$$

$$\xi_A^u = \frac{1}{t_\beta}, \quad \xi_A^{d,\ell} = -\frac{1}{t_\beta},$$

$$\text{Type-II: } \xi_h^u = s_{\beta-\alpha} + \frac{c_{\beta-\alpha}}{t_\beta}, \quad \xi_h^{d,\ell} = s_{\beta-\alpha} - c_{\beta-\alpha} t_\beta,$$

$$\xi_H^u = c_{\beta-\alpha} - \frac{s_{\beta-\alpha}}{t_\beta}, \quad \xi_H^{d,\ell} = c_{\beta-\alpha} + s_{\beta-\alpha} t_\beta,$$

$$\xi_A^u = \frac{1}{t_\beta}, \quad \xi_A^{d,\ell} = t_\beta,$$

$$a_h = s_{\beta-\alpha}, \quad a_H = c_{\beta-\alpha}.$$

Global signal fit to the
degenerate Higgs

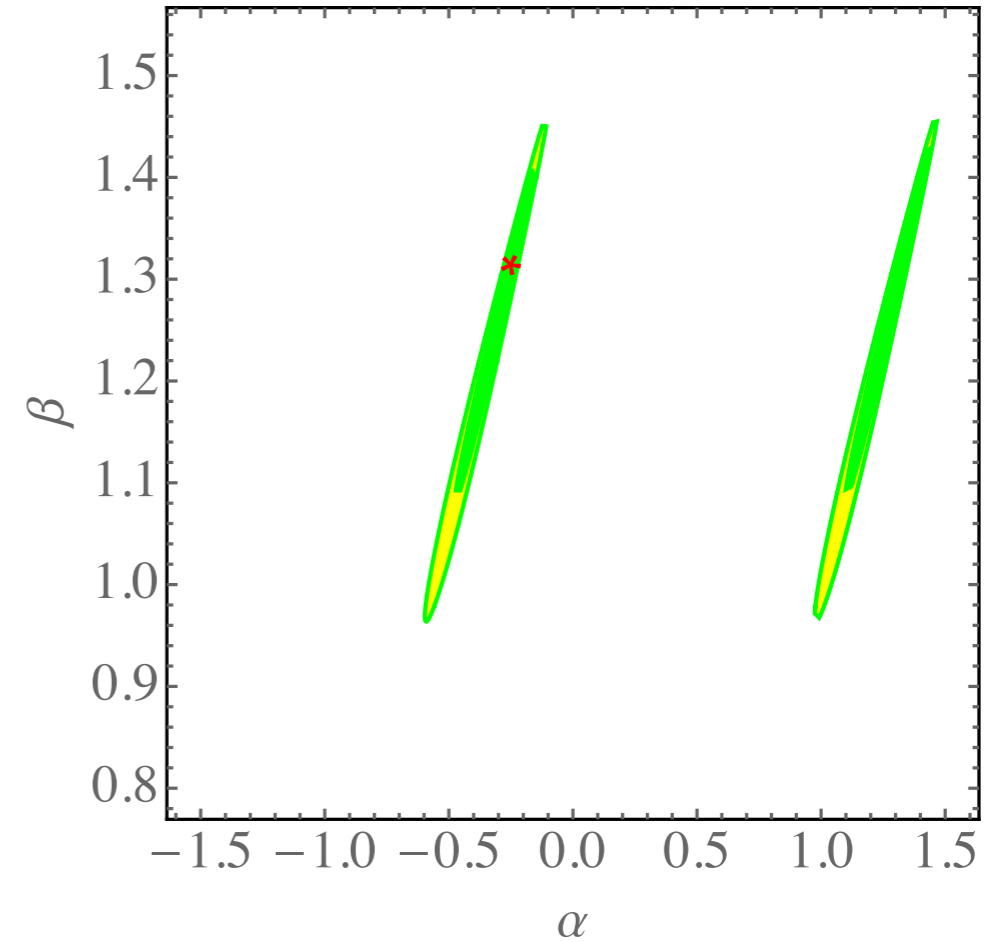
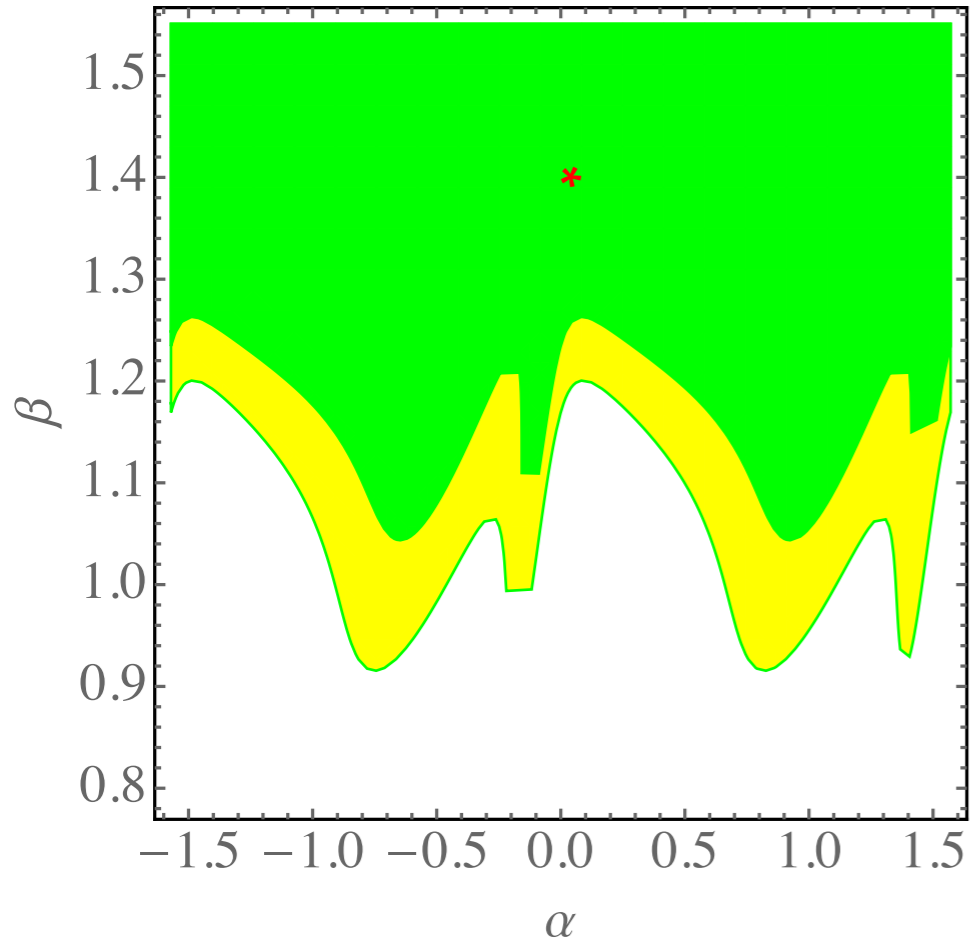
The Higgs signals @LHC run-1

Decays	Productions	ATLAS	Ref	CMS	Ref
$\gamma\gamma$	ggF	1.32 ± 0.38	[31]	$1.12^{+0.37}_{-0.32}$	[32]
$\gamma\gamma$	VBF	0.8 ± 0.7	[31]	$1.58^{+0.77}_{-0.68}$	[32]
$\gamma\gamma$	WH	1.0 ± 1.6	[31]	-	-
$\gamma\gamma$	ZH	$0.1^{+3.7}_{-0.1}$	[31]	-	-
$\gamma\gamma$	VH	-	-	$-0.16^{+1.16}_{-0.79}$	[32]
$\gamma\gamma$	ttH	$1.6^{+2.7}_{-1.3}$	[31]	$2.69^{+2.51}_{-1.81}$	[32]
ZZ	ggF,ttH,bbH	$1.7^{+0.5}_{-0.4}$	[33]	-	-
ZZ	ggF,ttH	-	-	$0.80^{+0.46}_{-0.36}$	[34]
ZZ	VBF,VH	$0.3^{+1.5}_{-0.9}$	[33]	$1.7^{+2.2}_{-2.1}$	[34]
W^+W^-	ggF	$1.02^{+0.29}_{-0.26}$	[35]	$0.74^{+0.22}_{-0.20}$	[36]
W^+W^-	VBF	$1.27^{+0.53}_{-0.45}$	[35]	$0.60^{+0.57}_{-0.46}$	[36]
W^+W^-	VH	-	-	$0.39^{+1.97}_{-1.87}$	[36]
$b\bar{b}$	ttH	1.5 ± 1.1	[37]	$1.2^{+1.6}_{-1.5}$	[38]
$b\bar{b}$	VH	$0.51^{+0.40}_{-0.37}$	[39]	1.0 ± 0.5	[40]
$\tau^+\tau^-$	ggF	$2.0^{+1.5}_{-1.2}$	[41]	1.07 ± 0.46	[42]
$\tau^+\tau^-$	VBF,VH	$1.24^{+0.59}_{-0.54}$	[41]	-	-
$\tau^+\tau^-$	VBF	-	-	0.94 ± 0.41	[42]
$\tau^+\tau^-$	VH	-	-	-0.33 ± 1.02	[42]

The Higgs signals @LHC run-II

Decays	Productions	ATLAS	Ref	CMS	Ref
$\gamma\gamma$	ggF	$0.80^{+0.19}_{-0.18}$	[43, 44]	$1.11^{+0.19}_{-0.18}$	[45]
$\gamma\gamma$	VBF	2.1 ± 0.6	[43, 44]	$0.5^{+0.6}_{-0.5}$	[45]
$\gamma\gamma$	VH	$0.7^{+0.9}_{-0.8}$	[43, 44]	$2.3^{+1.1}_{-1.0}$	[45]
$\gamma\gamma$	ttH	0.5 ± 0.6	[43, 44]	$2.2^{+0.9}_{-0.8}$	[45]
ZZ	ggF	$1.11^{+0.25}_{-0.22}$	[44, 46]	$1.20^{+0.22}_{-0.21}$	[47]
ZZ	VBF	$4.0^{+1.8}_{-1.5}$	[44, 46]	$0.05^{+1.03}_{-0.05}$	[47]
ZZ	VH	0 ± 1.9	[44, 46]	0 ± 2.83 , or 0 ± 2.66	[47]
ZZ	ttH	0 ± 3.9	[44, 46]	0 ± 1.19	[47]
W^+W^-	ggF	-	-	1.02 ± 0.27	[48]
W^+W^-	VBF	$1.7^{+1.2}_{-0.9}$	[49]	-	-
W^+W^-	WH	$3.2^{+4.4}_{-4.2}$	[49]	-	-
W^+W^-	VBF+VH	-	-	0.89 ± 0.67	[48]
$b\bar{b}$	VH	$1.20^{+0.42}_{-0.36}$	[50]	-	-
$\tau^+\tau^-$	ggF	-	-	0.84 ± 0.89	[51]
$\tau^+\tau^-$	VBF	-	-	$1.11^{+0.34}_{-0.35}$	[51]
$\tau^+\tau^-$	ttH	-	-	$0.72^{+0.62}_{-0.53}$	[52]

Global signal fit



$$\chi^2 = \sum_{\text{PD}} \left(\frac{\mu_{\text{th}}^{\text{PD}} - \mu_{\text{exp}}^{\text{PD}}}{\sigma_{\text{exp}}^{\text{PD}}} \right)^2,$$

$$\mu[XX \rightarrow h/H \rightarrow YY] = \frac{|\kappa_{hXX}\kappa_{hYY}|^2}{\Gamma_h/\Gamma_h^{\text{SM}}} + \frac{|\kappa_{HXX}\kappa_{HYY}|^2}{\Gamma_H/\Gamma_h^{\text{SM}}}.$$

$$\Delta M \equiv M_H - M_h \gg \Gamma_H + \Gamma_h$$

Best-fit points

$M_h \approx M_H$	Type-I	Type-II
$(c_{\beta-\alpha}, t_\beta, \Gamma_{\text{tot}})$	(0.21, 5.8, 4.2 MeV)	(0.01, 3.8, 41.4 MeV)
(Γ_h, Γ_H)	(4.11 MeV, 0.05 MeV)	(3.89 MeV, 37.49 MeV)
$(\text{Br}[h \rightarrow bb], \text{Br}[H \rightarrow bb])$	(58.80%, 8.68%)	(56.23%, 89.85%)
$(\text{Br}[h \rightarrow \tau^+\tau^-], \text{Br}[H \rightarrow \tau^+\tau^-])$	(6.44%, 0.95%)	(6.16%, 9.84%)
$(\text{Br}[h \rightarrow W^+W^-], \text{Br}[H \rightarrow W^+W^-])$	(20.35%, 77.86%)	(22.50%, -)
$(\text{Br}[h \rightarrow ZZ], \text{Br}[H \rightarrow ZZ])$	(2.50%, 9.56%)	(2.76%, -)
$(\text{Br}[h \rightarrow \gamma\gamma], \text{Br}[H \rightarrow \gamma\gamma])$	(0.21%, 1.22%)	(0.24%, -)
$(\text{Br}[h \rightarrow gg], \text{Br}[H \rightarrow gg])$	(8.73%, 1.29%)	(9.04%, 0.28%)

Table: The best-fit points of (α, β) for the mass-degenerate Higgs bosons of $M_h \approx M_H = 125$ GeV in both Type-I and Type-II 2HDM. The decay widths of (Γ_h, Γ_H) , and decay branching ratios are listed, where the decay branching ratios smaller than 10^{-4} are neglected.

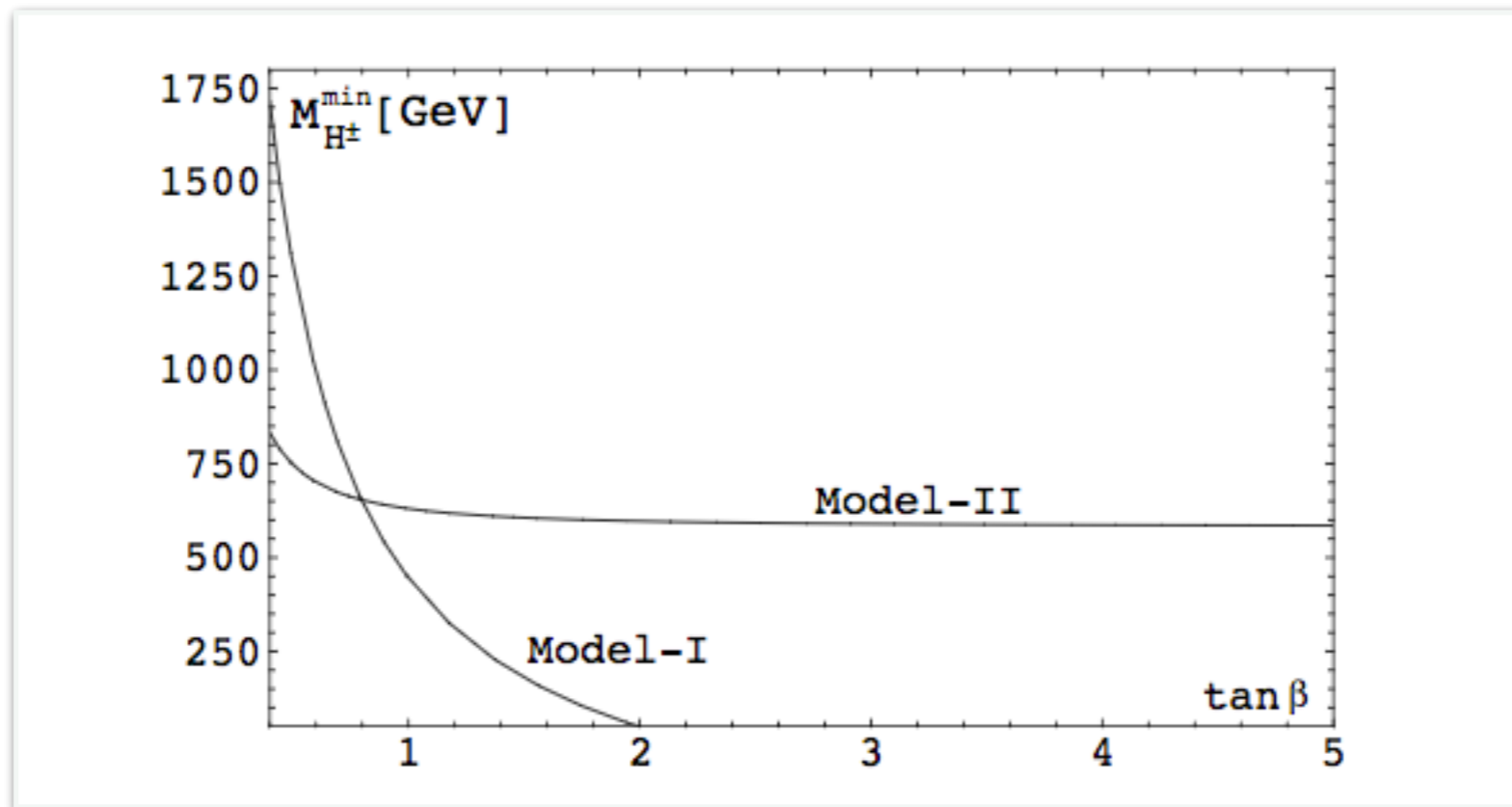
The constraints to the
2HDM with
degenerate Higgs

A set of constraints

- The charged Higgs boson mass bounds: from FCNC decays
- The unitarity & stability constraints to the 2HDM potential
- The EW precision tests: can be easily alleviated
- LHC direct searches for CP-odd scalar A

The charged Higgs boson

The Belle measurements of $b \rightarrow s + \gamma$ decay: 1608.02344



Misiak & Steinhäuser, 1702.04571

No constraints to M_{\pm} in Type-I, while $M_{\pm} \geq 590$ GeV in Type-II.

The unitarity && stability

- The perturbative unitarity constraints require the following quantities:

$$a_{\pm} = \frac{3}{2}(\lambda_1 + \lambda_2) \pm \frac{1}{2}\sqrt{9(\lambda_1 - \lambda_2)^2 + (2\lambda_3 + \lambda_4)^2},$$

$$b_{\pm} = \frac{1}{2}\left[(\lambda_1 + \lambda_2) \pm \sqrt{(\lambda_1 - \lambda_2)^2 + 4\lambda_4^2}\right],$$

$$c_{\pm} = \frac{1}{2}\left[(\lambda_1 + \lambda_2) \pm \sqrt{(\lambda_1 - \lambda_2)^2 + 4\lambda_5^2}\right],$$

$$f_+ = \lambda_3 + 2\lambda_4 + 3\lambda_5, \quad f_- = \lambda_3 + \lambda_5, \quad f_1 = f_2 = \lambda_3 + \lambda_4,$$

$$e_1 = \lambda_3 + 2\lambda_4 - 3\lambda_4, \quad e_2 = 2\lambda_3 - \lambda_5, \quad \rho_1 = \lambda_3 - \lambda_4,$$

to be $\leq 8\pi$.

- The stability constraints: $\lambda_{1,2} \geq 0$, $\lambda_3 \geq -\sqrt{\lambda_1\lambda_2}$, and $\lambda_3 + \lambda_4 - |\lambda_5| \geq -\sqrt{\lambda_1\lambda_2}$.

The unitarity && stability

- The quartic Higgs self couplings can be traded into physical inputs as:

$$\lambda_1 = \frac{M_h^2 s_\alpha^2 + M_H^2 c_\alpha^2 - m_{12}^2 t_\beta}{v^2 c_\beta^2},$$

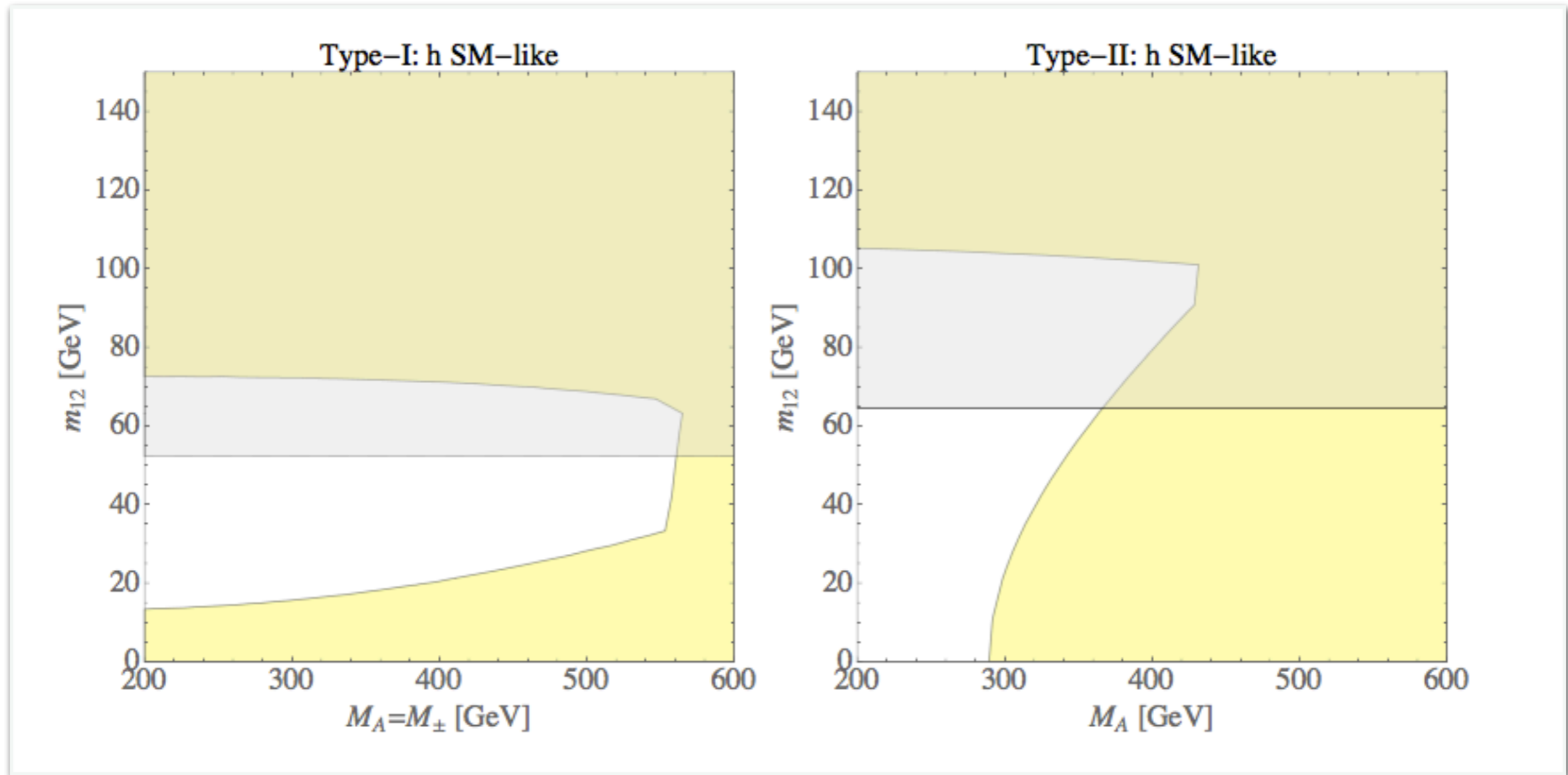
$$\lambda_2 = \frac{M_h^2 c_\alpha^2 + M_H^2 s_\alpha^2 - m_{12}^2 / t_\beta}{v^2 s_\beta^2},$$

$$\lambda_3 = \frac{1}{v^2} \left[\frac{(M_H^2 - M_h^2) s_\alpha c_\alpha}{s_\beta c_\beta} + 2M_\pm^2 - \frac{m_{12}^2}{s_\beta c_\beta} \right],$$

$$\lambda_4 = \frac{1}{v^2} \left(M_A^2 - 2M_\pm^2 + \frac{m_{12}^2}{s_\beta c_\beta} \right),$$

$$\lambda_5 = \frac{1}{v^2} \left(\frac{m_{12}^2}{s_\beta c_\beta} - M_A^2 \right).$$

The unitarity & stability



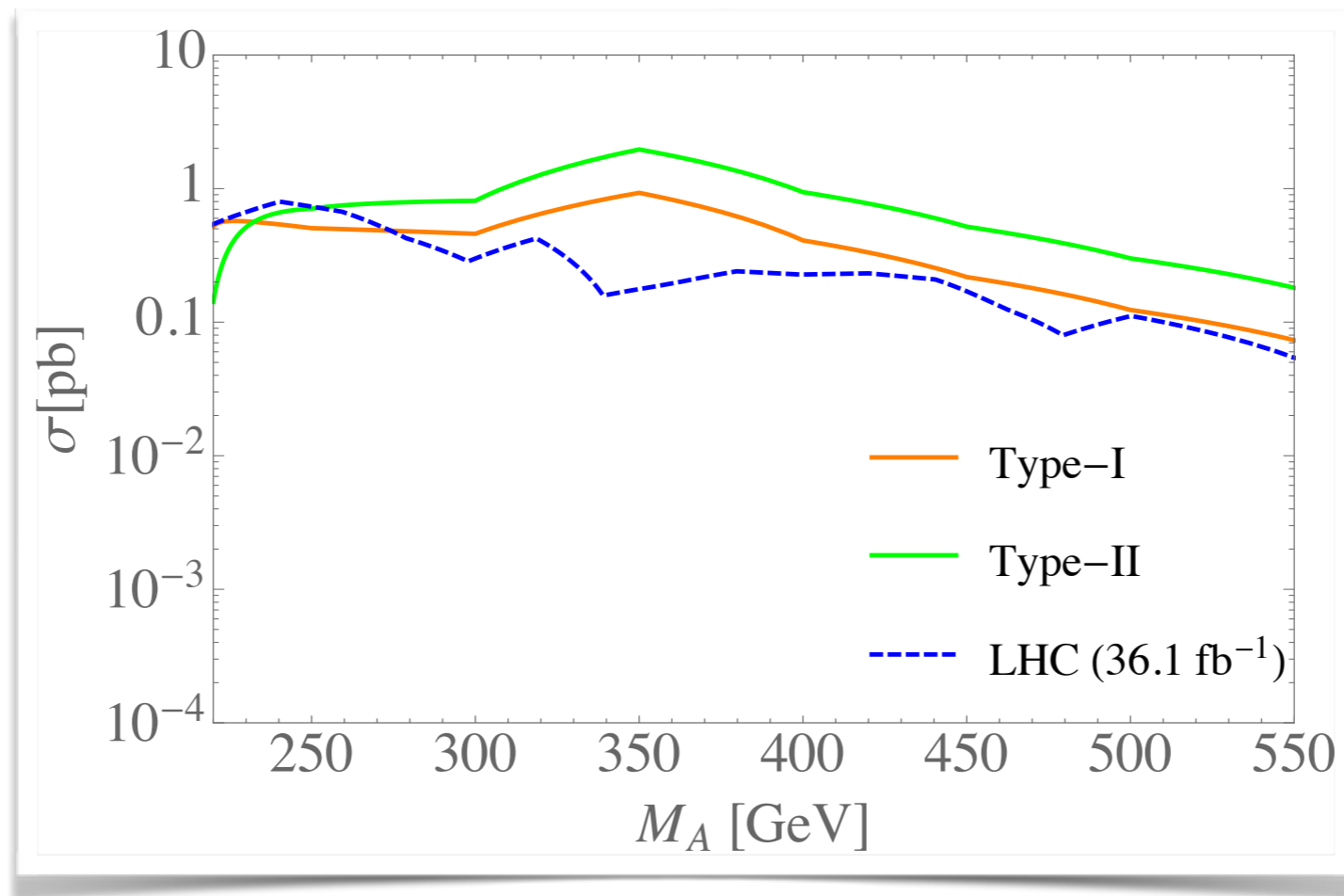
- Type-I : $20 \lesssim m_{12} \lesssim 50$ GeV, with $M_A = M_{\pm} \in (200, 280)$ GeV,
 Type-II : $0 \leq m_{12} \lesssim 60$ GeV,
 with $M_A \in (200, 250)$ GeV and $M_{\pm} = 600$ GeV.

Experimental tests of the degenerate Higgs

Indirect tests at LHC

- In the single $M_h \approx 125$ GeV 2HDM, the decay branching fraction of $A \rightarrow h + Z$ is suppressed by the alignment factor of $c_{\beta-\alpha} \sim \mathcal{O}(0.1) - \mathcal{O}(0.01)$.
- In the h/H mass-degenerate case, one has decay modes of $A \rightarrow h/H + Z$. The signal rates with the $b\bar{b} + \ell^+\ell^-$ final states are obtained as

$$\begin{aligned} \sigma_{\text{tot}} = & \sigma[pp \rightarrow AX] \times \left(\text{BR}[A \rightarrow hZ] \times \text{BR}[h \rightarrow b\bar{b}] \right. \\ & \left. + \text{BR}[A \rightarrow HZ] \times \text{BR}[H \rightarrow b\bar{b}] \right) \times \text{BR}[Z \rightarrow \ell^+\ell^-]. \end{aligned}$$

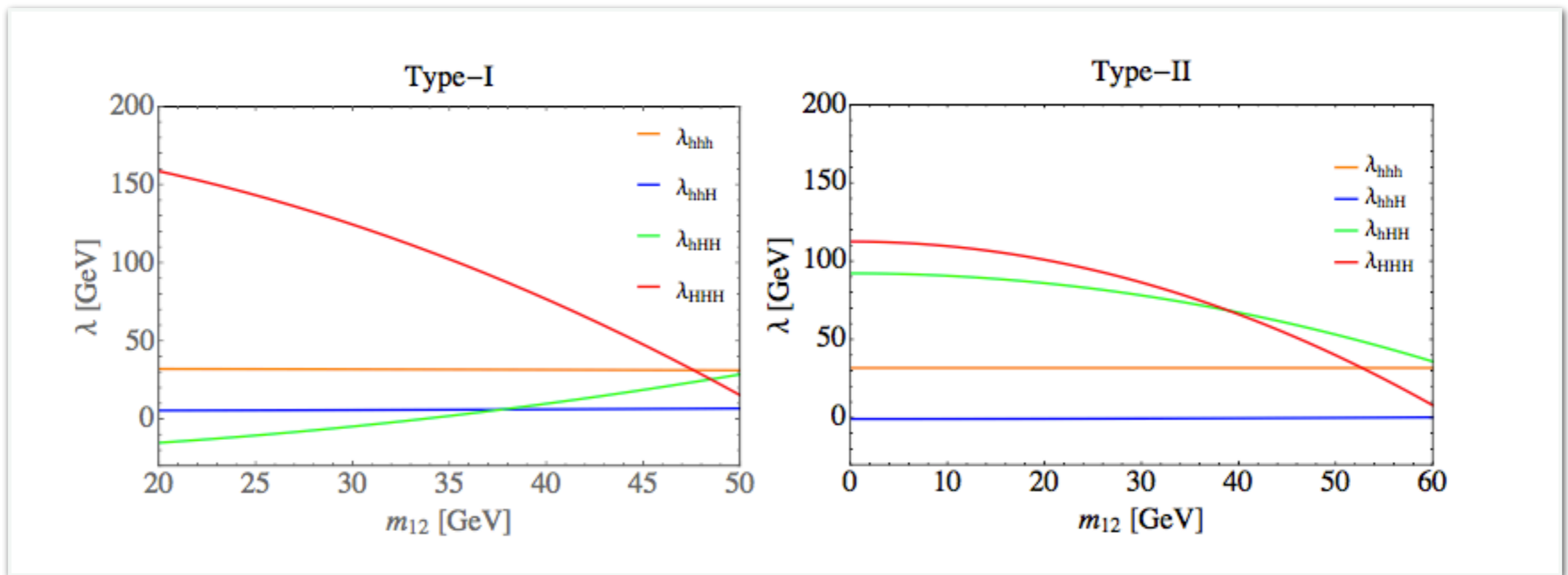


Direct tests: Higgs self-couplings & Higgs pairs

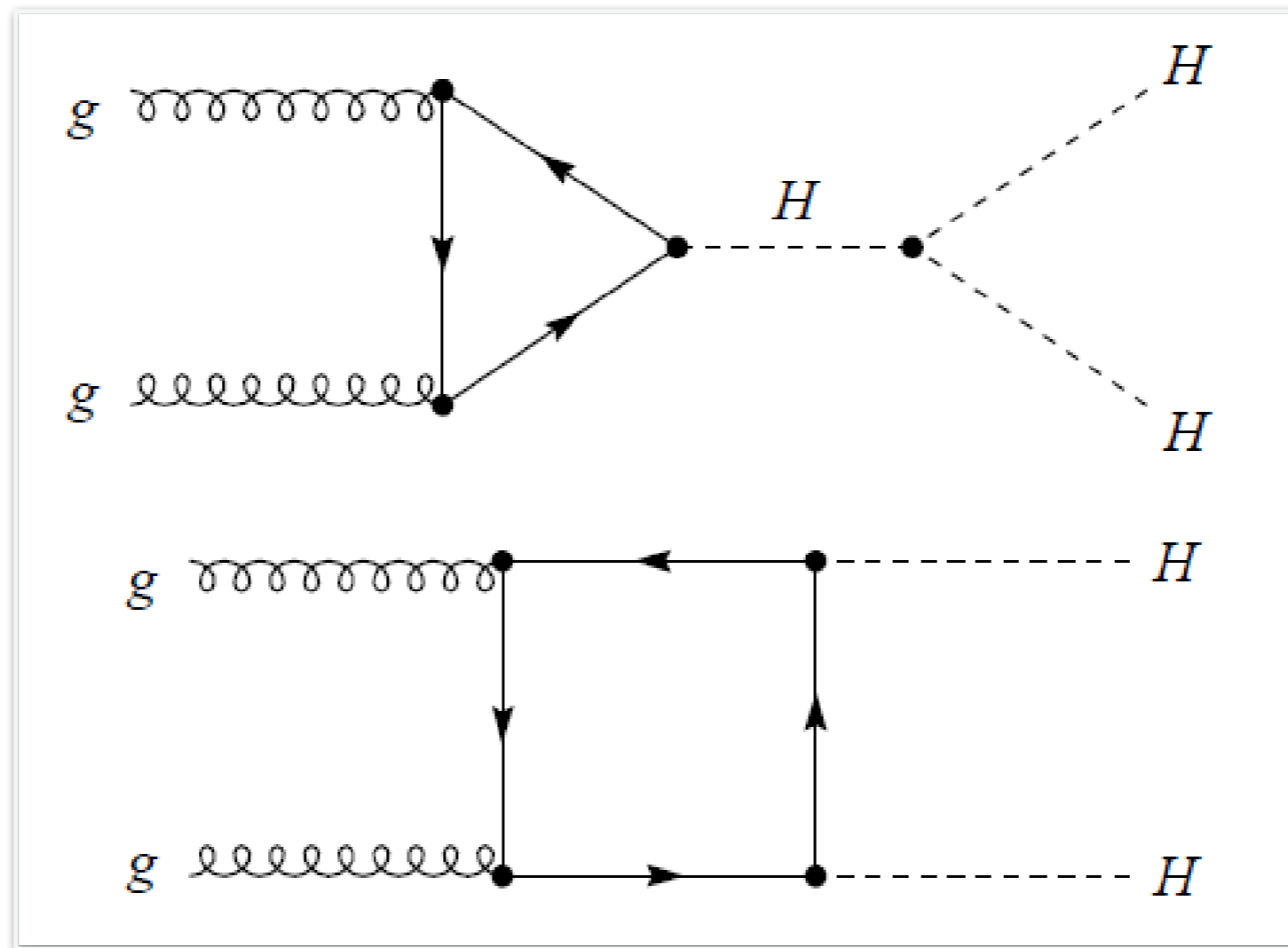
- The Higgs self-couplings can be measured through the Higgs pair productions at both LHC and future electron-positron colliders.
- There are more than one triple Higgs self-couplings involved for degenerate Higgs:

$$\begin{aligned} \lambda_{hhh} &= -\frac{1}{32vs_\beta^2c_\beta^2} \left[M_h^2(3s_{\alpha-\beta} + s_{3(\alpha-\beta)} - s_{3\alpha+\beta} - 3s_{\alpha+3\beta}) \right. \\ &\quad \left. + 4m_{12}^2(c_{3\alpha-\beta} + c_{\alpha-3\beta} + 2c_{\alpha+\beta}) \right], \\ \lambda_{hhH} &= \frac{c_{\alpha-\beta}}{2vs_\beta c_\beta} \left[(2M_h^2 + M_H^2)s_\alpha c_\alpha + m_{12}^2(1 - 3\frac{s_{2\alpha}}{s_{2\beta}}) \right], \\ \lambda_{hHH} &= \frac{s_{\beta-\alpha}}{2vs_\beta c_\beta} \left[-(M_h^2 + 2M_H^2)s_\alpha c_\alpha + m_{12}^2(1 + 3\frac{s_{2\alpha}}{s_{2\beta}}) \right], \\ \lambda_{HHH} &= -\frac{1}{32vs_\beta^2c_\beta^2} \left[M_H^2(c_{3(\alpha-\beta)} - c_{3\alpha+\beta} - 3c_{\alpha-\beta} + 3c_{\alpha+3\beta}) \right. \\ &\quad \left. + 4m_{12}^2(s_{\alpha-3\beta} - s_{3\alpha-\beta} + 2s_{\alpha+\beta}) \right]. \end{aligned}$$

Higgs self-couplings



Direct tests at the LHC: SM-like case



Direct tests at the LHC

- The relevant Higgs pair productions at the LHC are:

$$\sigma[gg \rightarrow hh], \quad \sigma[gg \rightarrow hH], \quad \sigma[gg \rightarrow HH]$$

- The parton-level cross section reads

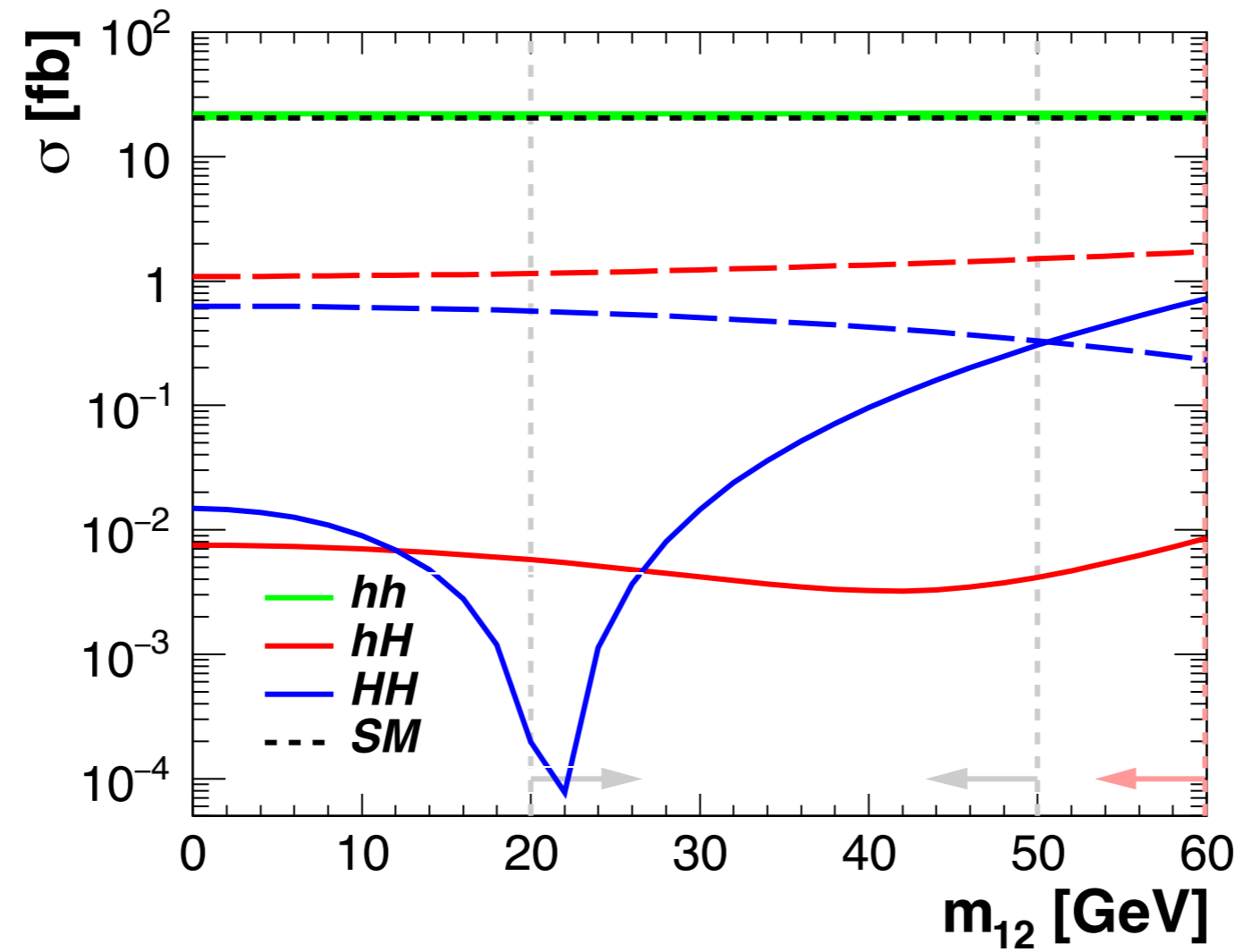
$$\frac{d\hat{\sigma}}{d\hat{t}} = K c^{ij} \frac{G_F^2 \alpha_s^2}{256(2\pi)^3} \left\{ \left| \sum_{q=t,b} (C_{\Delta}^{ij} F_{\Delta} + C_{\square}^{ij} F_{\square}) \right|^2 + \left| \sum_{q=t,b} C_{\square}^{ij} G_{\square} \right|^2 \right\}$$

with $K \approx 2$.

- The lab-frame cross section

$$\frac{d^2\sigma}{dM_{hh} dp_T} = \int_{\tau}^1 \frac{dx}{x} f_g(x, \mu_F) f_g\left(\frac{\tau}{x}, \mu_F\right) \frac{2M_{hh}}{s} \frac{d\hat{\sigma}}{dp_T}$$

Direct tests at the LHC



Direct tests at the LHC

- The signal rate estimations

$$\begin{aligned} \sigma[gg \rightarrow (hh, hH, HH) \rightarrow (XXYY)] &= \sigma[gg \rightarrow hh](\kappa_{XY}\text{Br}[h \rightarrow XX]\text{Br}[h \rightarrow YY]) \\ &+ \sigma[gg \rightarrow hH] \times (\dots) \\ &+ \sigma[gg \rightarrow HH](\kappa_{XY}\text{Br}[H \rightarrow XX]\text{Br}[H \rightarrow YY]) \end{aligned}$$

- The conservative estimations of significances for SM Higgs boson pairs at the HL-LHC:

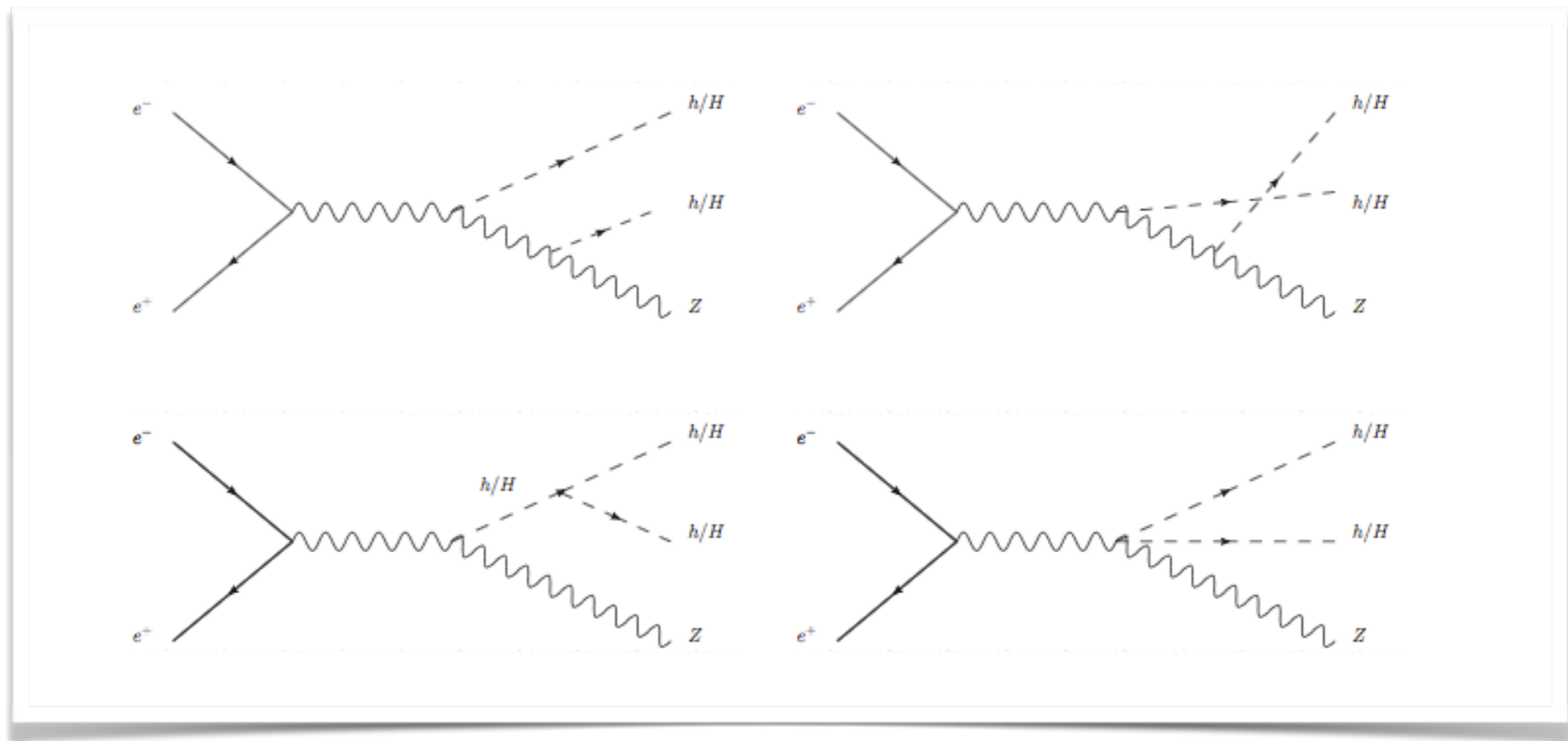
ATLAS : 1.05σ for $2b2\gamma$

CMS : $(0.39 \sigma, 1.6 \sigma)$ for $(4b, 2b2\gamma)$

- Not optimistic to probe the h/H mass-degenerate at the LHC:

	$4b$	$2b2\gamma$
Type-I (m_{12})	$\sim 0.44 \sigma$ (20 – 50 GeV)	1.61σ (20 GeV) 1.64σ (50 GeV)
Type-II (m_{12})	$\sim 0.43 \sigma$ (0 – 60 GeV)	1.71σ (0 GeV) 1.75σ (60 GeV)

Higgs pairs at the ILC



Higgs pairs at the ILC

- The cross sections at the ILC can be expressed as (Hikasa: Phys.Rev. D33 (1986) 3203)

$$\begin{aligned} \sigma = & \frac{1}{4} \left[(1 + P_{e^-})(1 + P_{e^+})\sigma_{RR} + (1 - P_{e^-})(1 - P_{e^+})\sigma_{LL} \right. \\ & \left. + (1 + P_{e^-})(1 - P_{e^+})\sigma_{RL} + (1 - P_{e^-})(1 + P_{e^+})\sigma_{LR} \right], \end{aligned}$$

where $\sigma_{LR} : (P_{e^+}, P_{e^-}) = (+1, -1)$ config.

- The ILC run at $\sqrt{s} = 500$ GeV with $\mathcal{L}_{\text{int}} = 4 \text{ ab}^{-1}$ shared by $(P_{e^+}, P_{e^-}) = (\pm 0.3, \mp 0.8)$:

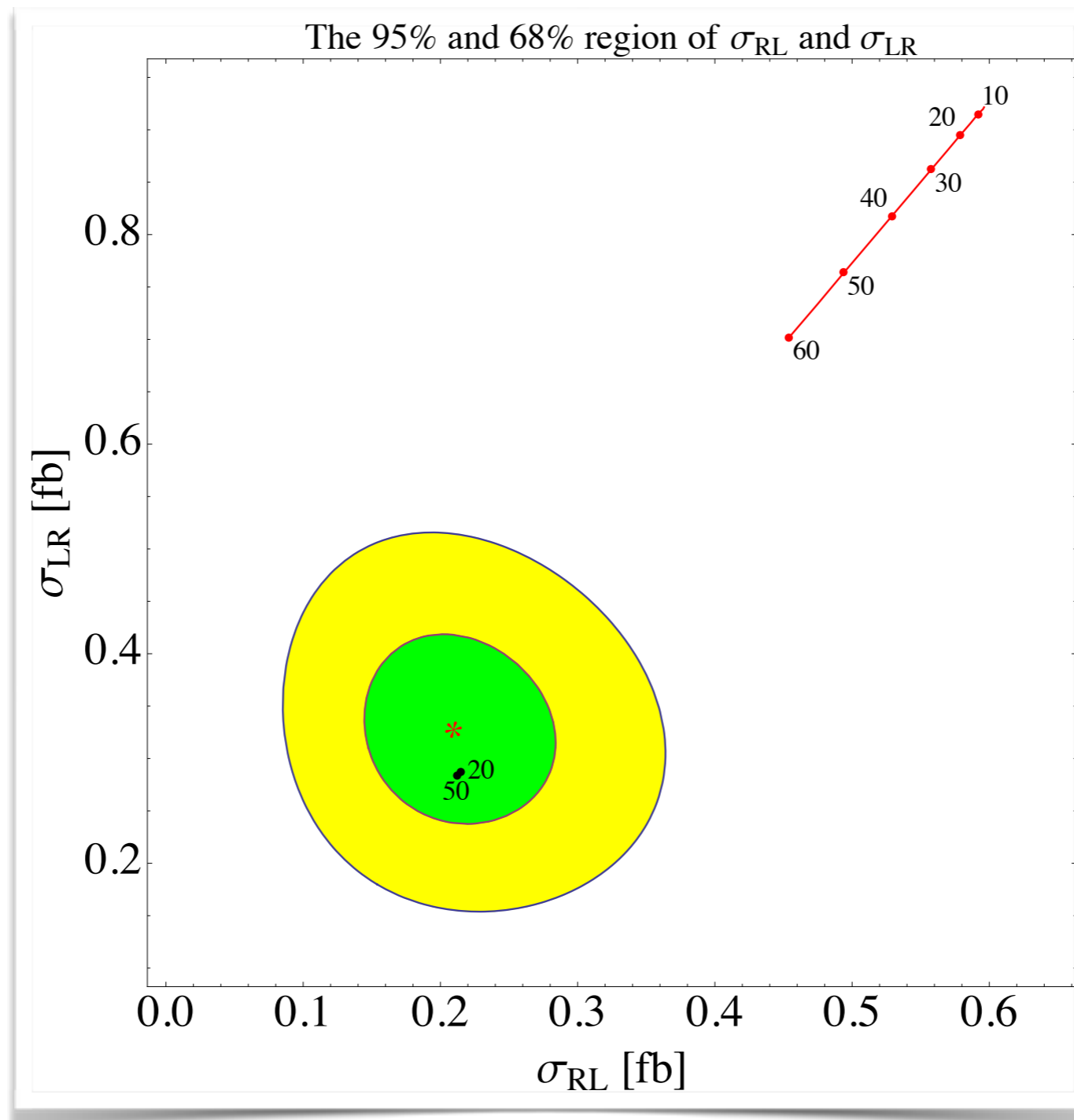
$$\begin{aligned} \sigma_{(+0.3, -0.8)} &= 0.585 \sigma_{LR} + 0.035 \sigma_{RL}, \\ \sigma_{(-0.3, +0.8)} &= 0.035 \sigma_{LR} + 0.585 \sigma_{RL}. \end{aligned}$$

Higgs pairs at the ILC

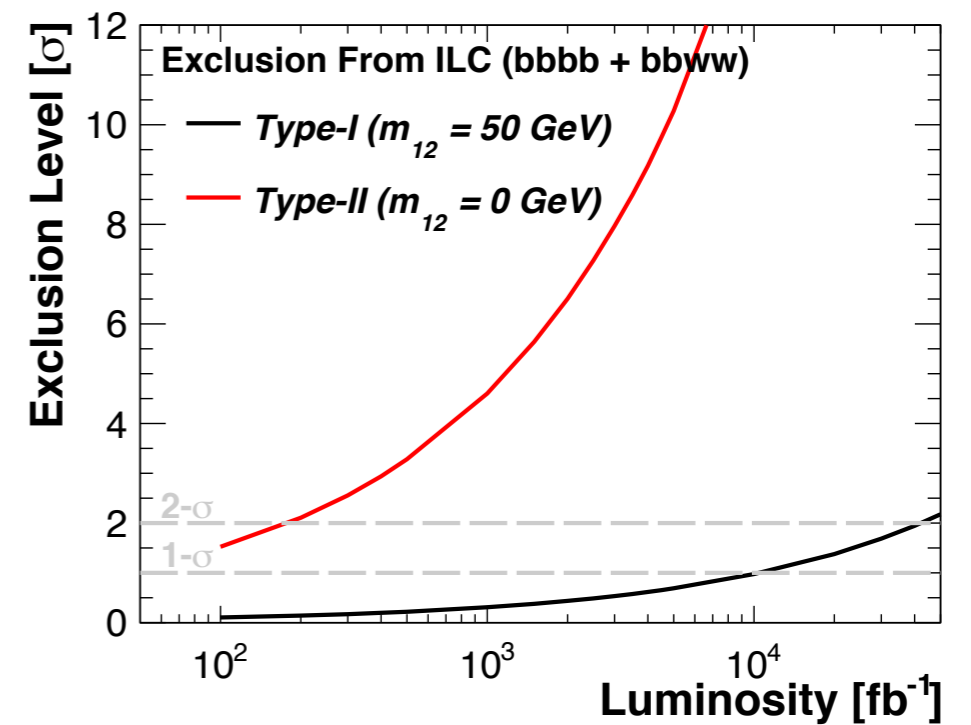
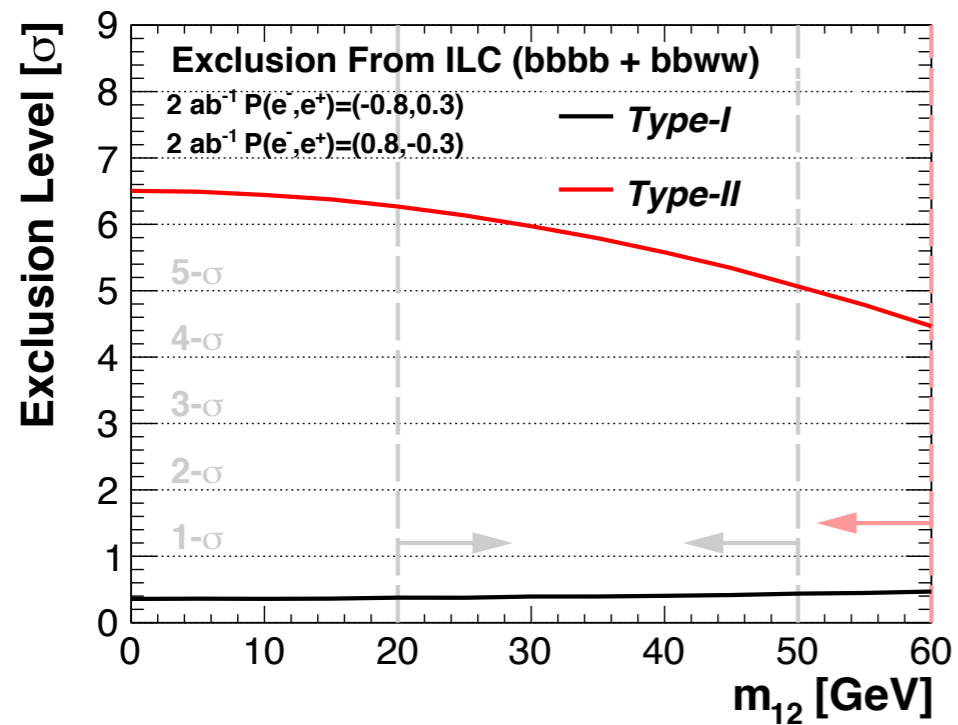
- The prospects (Claude Dürig: PhD thesis):

$P(e^+, e^-)$	Channel	Excess significance	Precision on σ_{ZH}
(0.3, -0.8)	$HH \rightarrow b\bar{b}b\bar{b}$	3.5σ	30.3%
(-0.3, 0.8)	$HH \rightarrow b\bar{b}b\bar{b}$	4.8σ	29.4%
(0.0, -0.8)	$HH \rightarrow b\bar{b}b\bar{b}$	3.5σ	34.7%
(0.0, 0.8)	$HH \rightarrow b\bar{b}b\bar{b}$	4.2σ	33.7%
(0.6, -0.8)	$HH \rightarrow b\bar{b}b\bar{b}$	4.2σ	28.7%
(-0.6, 0.8)	$HH \rightarrow b\bar{b}b\bar{b}$	5.5σ	27.8%
(0.3, -0.8)	$HH \rightarrow b\bar{b}W^+W^-$	1.91σ	–

Higgs pairs at the ILC



Higgs pairs at the ILC



Summary

- The future prospects of distinguishing the mass-degenerate Vs. the single resonance case, in the 2HDM.
- Direct measurements of Higgs signals: Higgs Yukawa && gauge couplings. We suggest to probe the Higgs self-couplings through Higgs pair productions.
- LHC: not likely to probe the scenario from the Higgs pairs, but can rule out this scenario by searching for A (very soon at run-II).
- Electron-positron colliders (ILC): direct measurements, can fully justify the Type-II case.

Thank you!